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Effects of the skew angle and road embankment length on the hydraulic performance of bridges on compound channels

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The flow pattern and bed shear stress are important factors to determine the scour potential regions in river bridges. This study has developed a 3D numerical model to simulate the flow around bridge abutments in a compound channel, and study the simultaneous effects of the bridge skew angle and contraction ratio (i.e. the roadway embankment length normal to the flow to the floodplain width) on the velocity distribution and the bed shear stress. After the model performance was verified, 13 cases were considered with different skew angles and contraction ratios. The results showed that the flow was more complex around the flow-splitting embankment than around the flow-guiding one, because of the flow-roadway embankment confrontation. At a zero skew angle, an increased contraction ratio increased the velocity around both abutments significantly; velocity increased by 67% and 40% when the contraction ratio rose from 0.25 to 0.50 and from 0.50 to 0.75, respectively. This velocity increase around the flow-splitting embankment was also visible for 15°, 30° and 45° angles; unexpectedly, the increase was less for the 45° angle than for the other cases. A shear stress study around the flow-guiding embankment showed that in all three cases an increased skew angle reduced the maximum shear stress mildly at contraction ratios of 0.25 and 0.50, and severely at 0.75. The trend of the maximum shear stress variations around the flow-splitting embankment was different for different contraction ratios – at a contraction ratio of 0.25 it was the highest at 45°, but for contraction ratios of 0.50 and 0.75, the maximum was 30°.

INTRODUCTION

Bridges are the most important and costly communication components in transportation systems. Since the most serious river-bridge destruction factor during floods is the scouring around its foundation and abutments (Aksoy & Eski 2016), knowing the scour-affecting factors such as the flow pattern and bed shear stress is vital.

A very important point about flows in natural waterways is the main channel overflow and water spreading into the surrounding lands known as the floodplain. Here, due to such specific geometrical and hydraulic conditions as low depth, high roughness and existing obstacles, the flow pattern becomes complicated and its analytical study becomes difficult. Under floodplain flow conditions, the high flow velocity difference between the main channel and the floodplain creates an interaction zone at the inter-subsection boundaries, and the velocity exchange in this region causes great energy loss and shear stress.

In studying the hydraulics of compound sections, the analysis becomes more complicated when the floodplain experiences a width variation for any reason. Effects of the floodplain cross-section variations will cause the flow to change from uniform to non-uniform and extra momentum/turbulence exchange between the main channel and the floodplain. One reason for this is constructing bridges in compound sections. At the bridge section, some part of the river floodplain is usually blocked for economic reasons (e.g. to reduce span), causing the river to face a contraction in its path during a flood. Figure 1 shows a schematic view of a bridge on a channel with a compound cross-section where the roadway embankment has blocked part of the floodplain with a skew angle θ . In this study, the ratio of the roadway embankment length normal to flow (L_R) to the floodplain width (B_f) is defined as the contraction ratio (L_R/B_f).

Numerous numerical and experimental studies have been done so far to determine

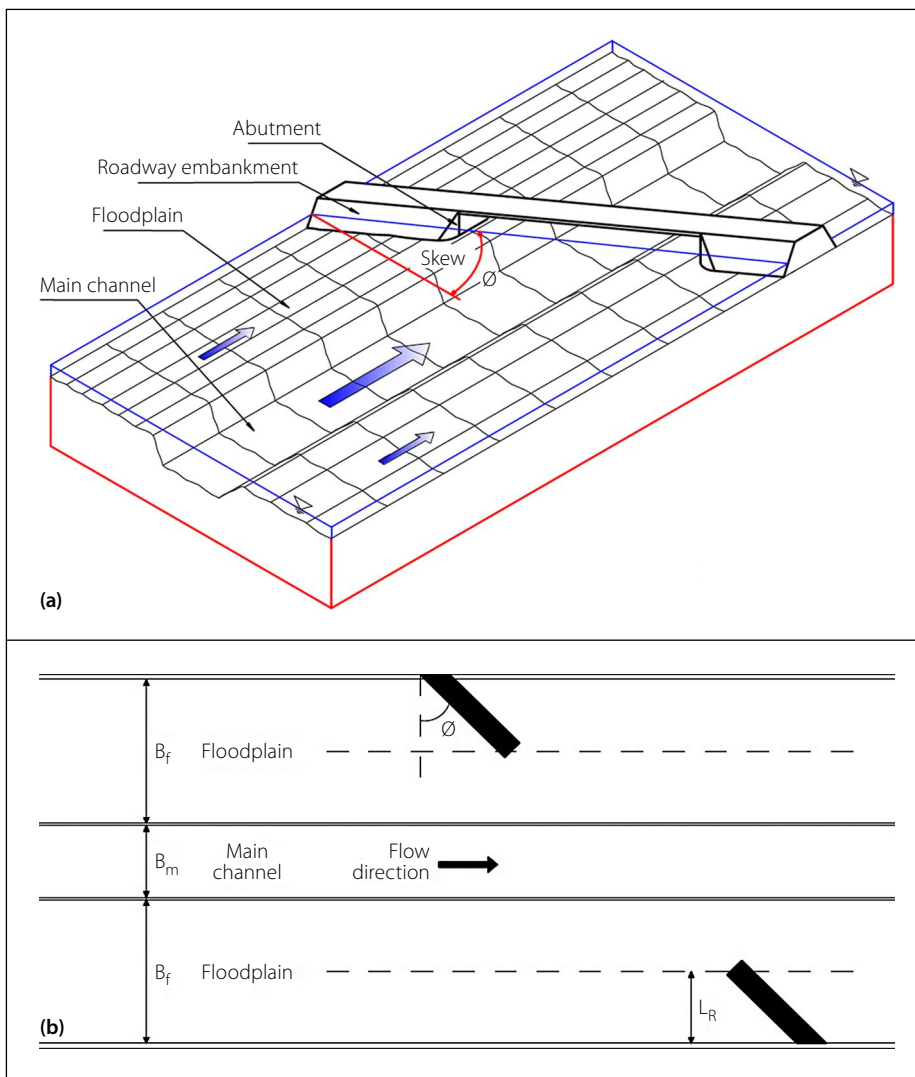


Figure 1 A schematic view of a skew bridge on a compound channel where part of the floodplain has been blocked by the road embankment: (a) 3D view, (b) Plan

the flow pattern and scouring around rectangular-section bridge abutments. Kandasamy (1989) did extensive studies on the effects of the abutment length on the clear water scouring in a rectangular channel and presented the results for four different zones. Dongol (1994) studied the effects of the abutment length on the scouring and plotted his data with those of Tey (1984) and Kandasamy (1989) for wing-shape abutments, and with those of Kwan (1984) for semicircular abutments. In a similar work, Melville (1997) added more data to those of Dongol (1994) for wing-shape abutments and abutments with vertical and inclined walls.

Molinas *et al* (1998) performed experiments on vertical abutments for flows with Froude numbers in the range 0.3–0.9 and contraction ratios 0.1, 0.2 and 0.3 (ratio of the abutment length normal to the flow to the total channel width) and showed that the bed shear stress and velocity at the abutment increased up to 10 times compared to those upstream. Ahmed and

Rajaratnam (2000) studied the flow field around a wing-shape abutment and concluded that the approaching flow turned into a complex 3D one upstream and around the abutment. They also found that the bed shear stress around the abutment (τ/τ_0) reached a maximum value equal to 3.63 near the nose (τ = bed shear stress and τ_0 = shear stress of the bed approaching flow). Teruzzi *et al* (2009) studied shear stresses on the bottom wall near the bridge abutment numerically, analysed the 3D flow against the abutment with emphasis on its effects on shear stresses and pressure gradients, and discussed their scouring potentials. Kara *et al* (2015) studied the effects of the free surface modelling on the flow around the bridge abutments using both the level-set and rigid-lid methods along with the LES (large eddy simulation) turbulence model and showed that the turbulence structure of this flow was strongly influenced by the water surface deformation, while the bed shear stress was quite similar for both models.

Various researchers have studied the flow pattern and scouring around bridge abutments in compound sections besides simple rectangular ones. By proposing a 2D depth-averaged turbulence model for a free surface flow in a 2-stage channel around bridge abutments on the floodplain, Biglari and Sturm (1998) concluded that the measured and computed depth-averaged velocities agreed well, but they emphasised that more 3D studies were needed to check the compound channel effects on the flow characteristics in regions close to obstacles. Sturm and Chrysochoides (1998) experimentally studied scouring around various lengths of bridge abutments under clear water conditions in a compound channel. To predict the scouring depth, Sturm (1999) proposed an equation that depended on local values of hydraulic variables near the nose of the abutment. In another study, Sturm (2006) did some tests on different compound-section geometries considering the main channel–floodplain flow distribution in the contracted section, sediment size, backwater caused by the abutment obstruction, abutment shape, velocity of the approaching flow, and the flow depth. Developing a 2D depth-integrated model and a laboratory hydraulic version, Morales and Ettema (2013) studied the velocity and shear stress in the flow fields around bridge abutments in compound channels, considering different discharges, channel/abutment geometrical parameters and flow patterns. Using the existing Flow-3D commercial model, Kocaman (2014) evaluated the computational fluid dynamics performance in predicting the water surface level profile at bridges in compound channels with and without piers; the studied bridges were single- and/or double-span with semicircular openings, and single-span with semi-ellipsoid opening. He used the standard $k-\epsilon$ turbulence model to estimate the eddy viscosity, compared the numerical and laboratory results, and showed that they conformed well to those of the afflux profile upstream of the bridge. In another study, Atabay *et al* (2017) proposed an analytical method to calculate the height of afflux around bridges in compound channels. Chua *et al* (2019) used the LES method to investigate the flow/turbulence structure around varying-length bridge abutments in a channel with an asymmetric compound section.

Another important point in bridge hydraulics is the skew angle and its impact on the flow and bridge-pier scouring. This

has been studied by many researchers so far, e.g. Lança *et al* (2012), Yang *et al* (2017), and Shahhosseini and Yu (2019). Studies on the effects of the bridge skew angle on abutments are more limited. Using a 3D numerical model and three groyne angles (45°, 90° and 135°) relative to the upstream flow, Haltigin *et al* (2007) studied the scouring-hole geometry near the groyne and showed that different lengths/angles of these structures affected it; the highest and lowest erosion belonged to 45° and 135°, respectively. Seckin (2007) performed some tests to study the abutment skew angle effects on the afflux in bridges in compound channels. Using the Flow-3D commercial model, Erduran *et al* (2012) simulated the flow around bridge skew angles in compound channels and compared the water surface profile with the test results obtained at the Birmingham University Hydraulic Lab; the comparison showed good agreement.

As mentioned earlier, numerical/experimental studies done, so far, to investigate the effects of the contraction ratio on the flow/scouring pattern in rectangular/compound channels are numerous, but those on the bridge skew angle are few. This study has utilised a 3D numerical model to investigate the simultaneous effects of the bridge skew angle and contraction ratio in compound channels on the flow patterns and the bed shear stress.

Setting up of 3D model

To simulate the flow, the 3D Navier Stokes equations and continuity equation have been used. By considering the Boussinesq approximation with incompressible flow conditions, these will be as follows (Schlichting 1979):

$$\frac{\partial u_j}{\partial x_j} = 0 \quad (1)$$

$$\frac{\partial u_i}{\partial t} + \frac{\partial (u_i u_j)}{\partial x_j} = -\frac{1}{\rho} \frac{\partial P}{\partial x_i} + \frac{\partial}{\partial x_j} \left[(\nu_t) \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \right] + \frac{f_i}{\rho} \quad (2)$$

where t , g , P , ρ , ν_t , f_i and u_i are time, gravitational acceleration, pressure, fluid density, eddy viscosity, body force and the velocity Cartesian components in the x_i directions, respectively. In the present study, the horizontal eddy viscosity (ν_t^h) is determined by the Smagorinsky model, while the vertical eddy viscosity (ν_t^v) is computed by the 1D k - ϵ model. The horizontal eddy viscosity is specified according to the Smagorinsky model by the following equation (Smagorinsky 1963).

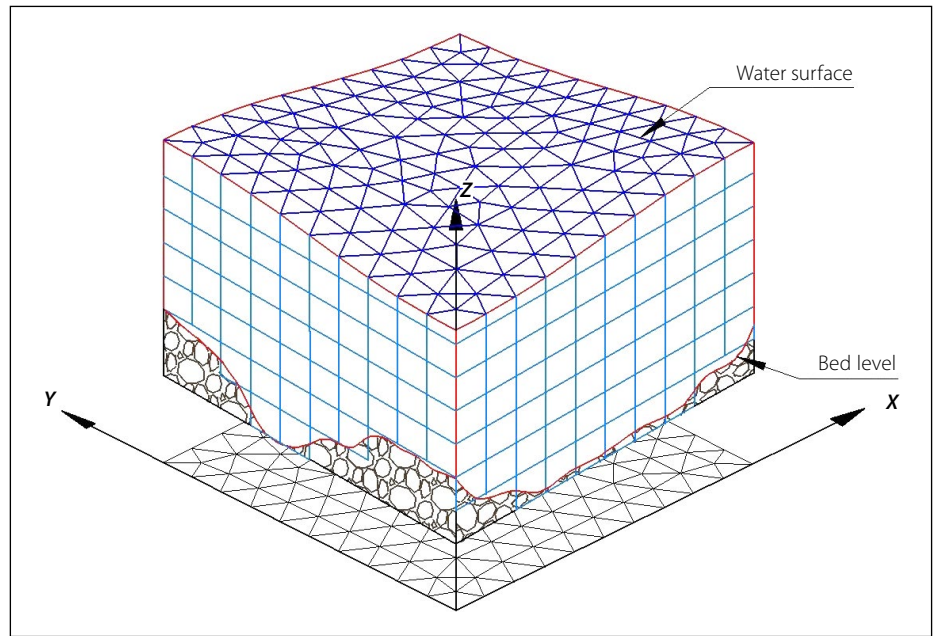


Figure 2 Three-dimensional grid layout – unstructured (triangular) and structured grids in the horizontal and vertical planes, respectively

$$\nu_t^h = (C_s \Delta)^2 \sqrt{2 \cdot \overline{S_{ij}} \cdot \overline{S_{ij}}} \quad (3)$$

where C_s is the Smagorinsky constant varying from 0.1 to 0.8, Δ is the length scale for the grid filter and $\overline{S_{ij}}$ is the filtered strain-rate tensor $\left[\overline{S_{ij}} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \right]$.

The 1D k - ϵ model uses transport equation for turbulent kinetic energy, k , and dissipation of turbulent kinetic energy, ϵ , in the vertical direction. The vertical eddy viscosity from the 1D k - ϵ model is as follows:

$$\nu_t^v = C_\mu \frac{k^2}{\epsilon} \quad (4)$$

In the above expression, C_μ is an empirical coefficient, which is usually a constant equal to 0.09 and k and ϵ are computed from the following transport equations:

$$\frac{\partial k}{\partial t} = \frac{\partial}{\partial z} \left(\nu_t^v \frac{\partial k}{\partial z} \right) + P - \epsilon \quad (5)$$

$$\frac{\partial \epsilon}{\partial t} = \frac{\partial}{\partial z} \left(\nu_t^v \frac{\partial \epsilon}{\partial z} \right) + C_{1\epsilon} \frac{\epsilon}{k} P - C_{2\epsilon} \frac{\epsilon^2}{k} \quad (6)$$

where $P = \nu_t^v \left[\left(\frac{\partial u}{\partial z} \right)^2 + \left(\frac{\partial v}{\partial z} \right)^2 \right]$ is the

production term due to velocity shear; σ_k , σ_ϵ , $C_{1\epsilon}$ and $C_{2\epsilon}$ are empirical parameters with constant values equal to 1, 1.3, 1.44 and 1.92 respectively (MIKE3 2007).

In regions of particular interest, small mesh-size grids will lead to more precise results and the computational domain should have a longer length to either connect the region to the closest known

boundary (which may be quite far from the river bridge) or prevent unwanted boundary reflections near the model region; fine grids cause inefficient computations in long regions. One solution to this problem is to use unstructured meshes that form finer grids near the bridge and coarser ones elsewhere in the region; here, an unstructured triangular mesh has been used to discretize the domain in the horizontal direction, and a structured grid has been used to divide it in the vertical direction to form prisms with triangular horizontal faces with similar x and y for each corresponding vertex. Such grids (Figure 2) fit into complicated geometries and enable local mesh refinements.

The finite volume fractional step method was used to solve the governing equations. The transport equation (advection and diffusion) is solved in the horizontal plane using the second-order explicit scheme while the Crank-Nicholson implicit scheme is deployed in the vertical direction. In the developed 3D model, a hydrostatic pressure assumption is applied at top layer cells. More details about boundary conditions, numerical scheme, free surface modelling and solving the system of equations can be found in Namin *et al* (2004) and Kilanehei *et al* (2011).

MODEL VALIDATION

To evaluate the performance of the developed numerical model, its results were compared with those of the laboratory tests used on unsymmetrical compound channels to investigate the velocity and shear stress parameters. To assess vortexes, the flow around

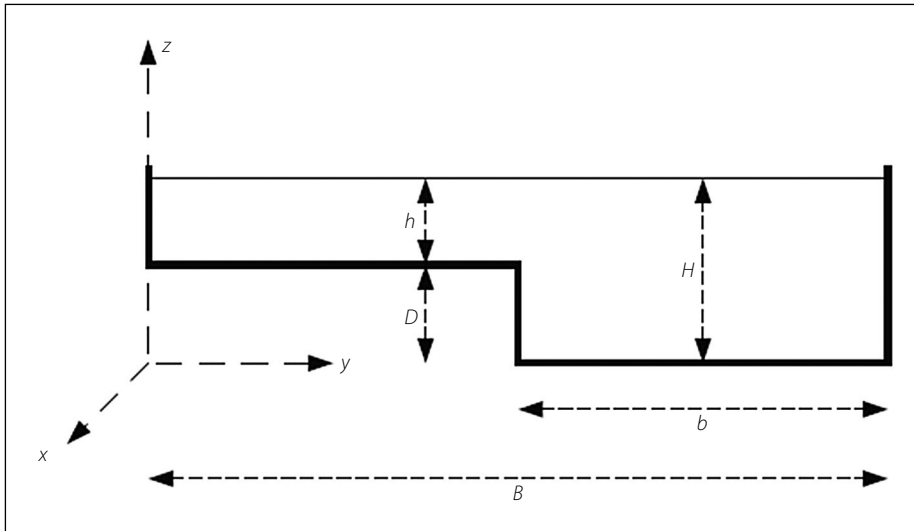


Figure 3 Geometrical configuration of the compound cross channel (Tominaga & Nezu 1991)

impermeable groynes were used on one side of the symmetrical compound channel.

Unsymmetrical compound channel

To validate the model for flow simulation capability in compound channels, the

laboratory data of Tominaga and Nezu (1991) were used; they did their tests using Laser Doppler Anemometry (LDA). In the laboratory model, the main channel width (b) was 20 cm and the water depth (H) was 8 cm, and there

was a floodplain with a width ($B-b$) equal to 20 cm only on one side (Figure 3). The water depth (h) in the floodplain was 4 cm and the 8 376 cm³/sec discharge was simulated in a flume 12.5 m long with a slope of 0.0006397.

These conditions were also applied to the numerical model. Figure 4 compares the longitudinal velocity distribution along different transverse sections. In this figure, U is the x -direction velocity, U_{max} is the maximum x -direction velocity, Y is the y -direction distance from the floodplain wall and Z is the z -direction distance from the main channel bed. The solid-line graph is related to the current numerical study and the dotted one corresponds to the laboratory model. When $Y < 20$, the section

Table 1 RMSE and MAPE values of transverse sections

	Y/H=1.2	Y/H=1.8	Y/H=2.75	Y/H=3.0	Y/H=3.65	Y/H=4.5	Y/H=4.75
RMSE (%)	1.98	1.96	2.88	2.77	2.49	2.06	3.66
MAPE (%)	2.38	1.87	2.83	2.75	2.65	2.09	4.25

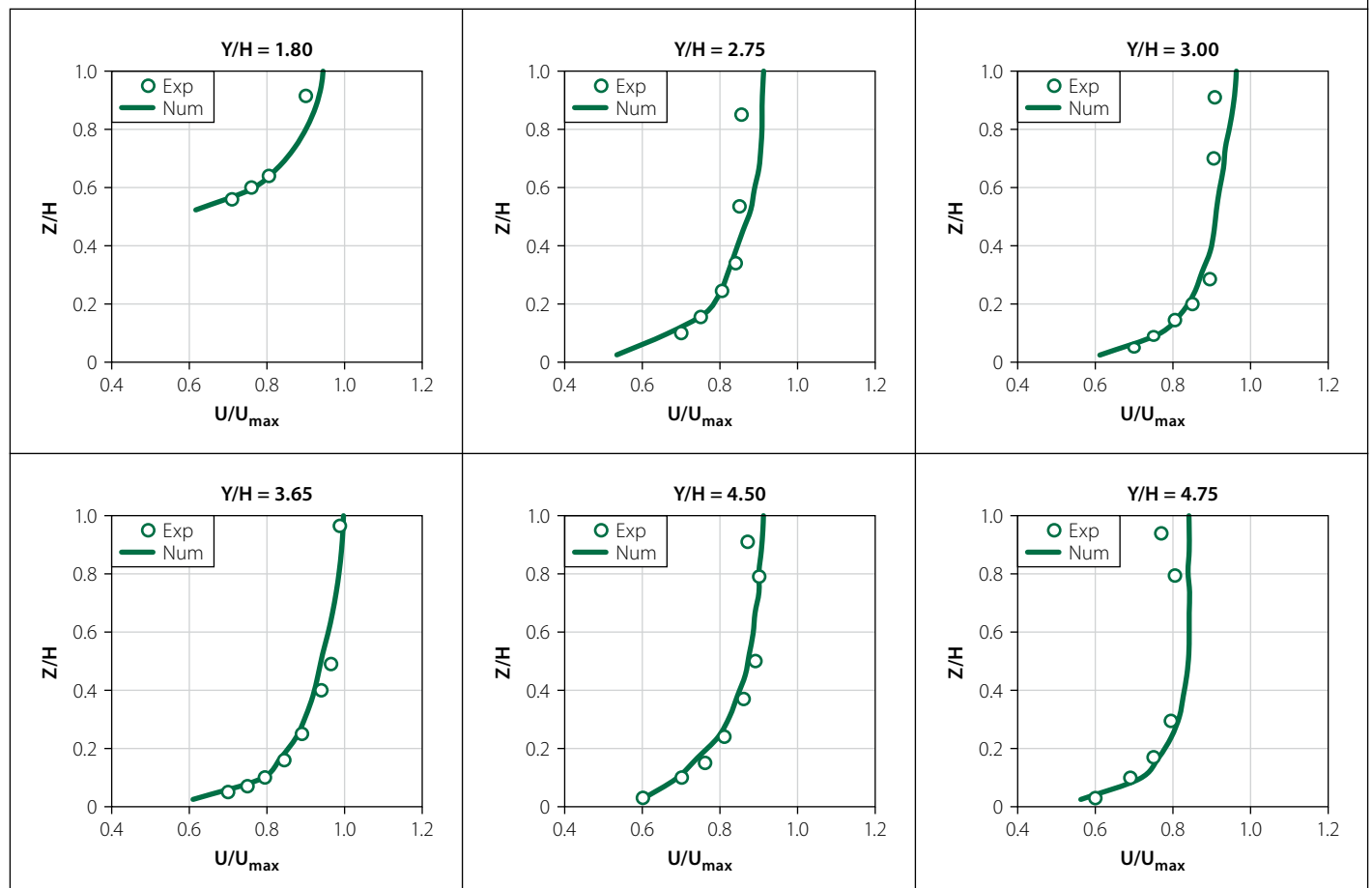


Figure 4 Longitudinal velocity distribution along different transverse sections

lies in the floodplain and when $Y > 20$, it lies in the main channel.

To analyse the results quantitatively, the root mean square error (RMSE) and the mean absolute percentage error (MAPE) were determined as follows:

$$RMSE = \sqrt{\frac{\sum_{i=1}^n (U_i^{Exp} - U_i^{Num})^2}{n}} \times 100 \quad (7)$$

$$MAPE = \left[\frac{1}{n} \sum_{i=1}^n \left| \frac{U_i^{Exp} - U_i^{Num}}{U_i^{Exp}} \right| \right] \times 100 \quad (8)$$

where *Num* and *Exp* superscripts represent the numerical and experimental model results, respectively. The RMSE and MAPE values are presented in Table 1 (page 47) for each transverse section. The mentioned values show the accuracy and consistency of numerical prediction. It is worth mentioning that the slight error increase in the numerical calculations at transverse sections near the walls ($Y/H = 2.75$, $Y/H = 3.0$ and $Y/H = 4.75$) could be due to the secondary flows.

As mentioned before, the bed shear stress distribution is important in sediment transport and scouring studies. Bed shear stress distribution for laboratory and current studies are shown in Figure 5; the dots relate to the former and the solid line corresponds to the latter (here τ_w is the average shear stress).

As shown, when the floodplain approaches the main channel junction, the shear stress increases due to their momentum transfer. RMSE and MAE are 9.09% and 7.69%, respectively, showing the accuracy and consistency of the numerical predictions.

Flow around impermeable groynes on one side of the symmetrical compound channel

The laboratory data of Ahmed *et al* (2010) was used to assess the numerical model's flow simulation capability in compound channels with obstacles on its floodplain. The experiments, conducted at Saitama University in Japan, used a 15 m long, 0.5 m deep, 0.5 m wide symmetrical compound section flume with a main channel width (B) of 0.1 m and a total water depth (H) of 0.24 m, and two symmetrical floodplains with width b equal to 0.2 m (floodplain relative width $b/B = 2.0$). The flow was a steady type with a discharge (Q) equal to 0.015 m³/s and a floodplain water depth (h) equal to 0.08 m (Figure 6).

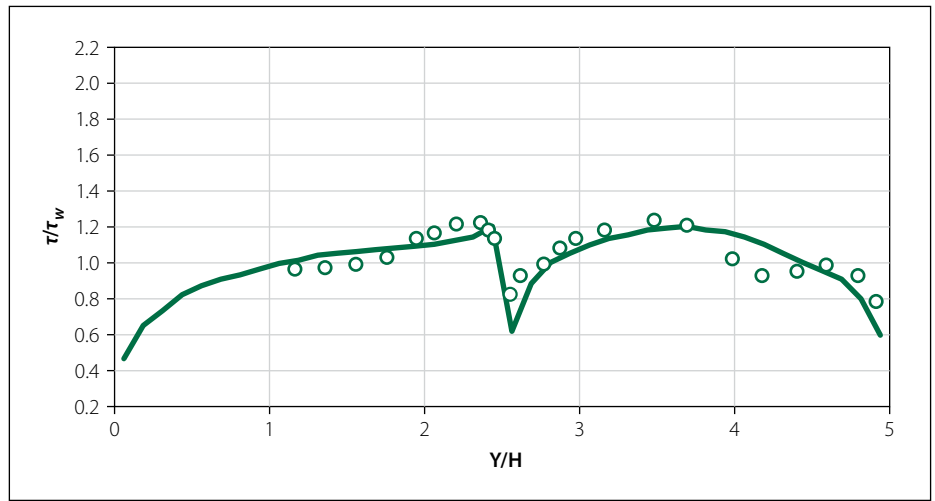


Figure 5 Bed shear stress distribution in the compound channel

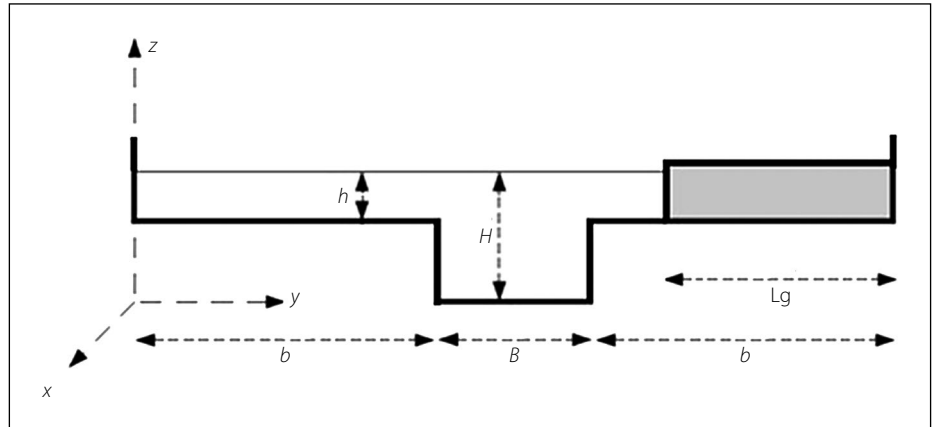


Figure 6 Geometrical configuration of the compound channel cross section (Ahmed *et al* 2010)

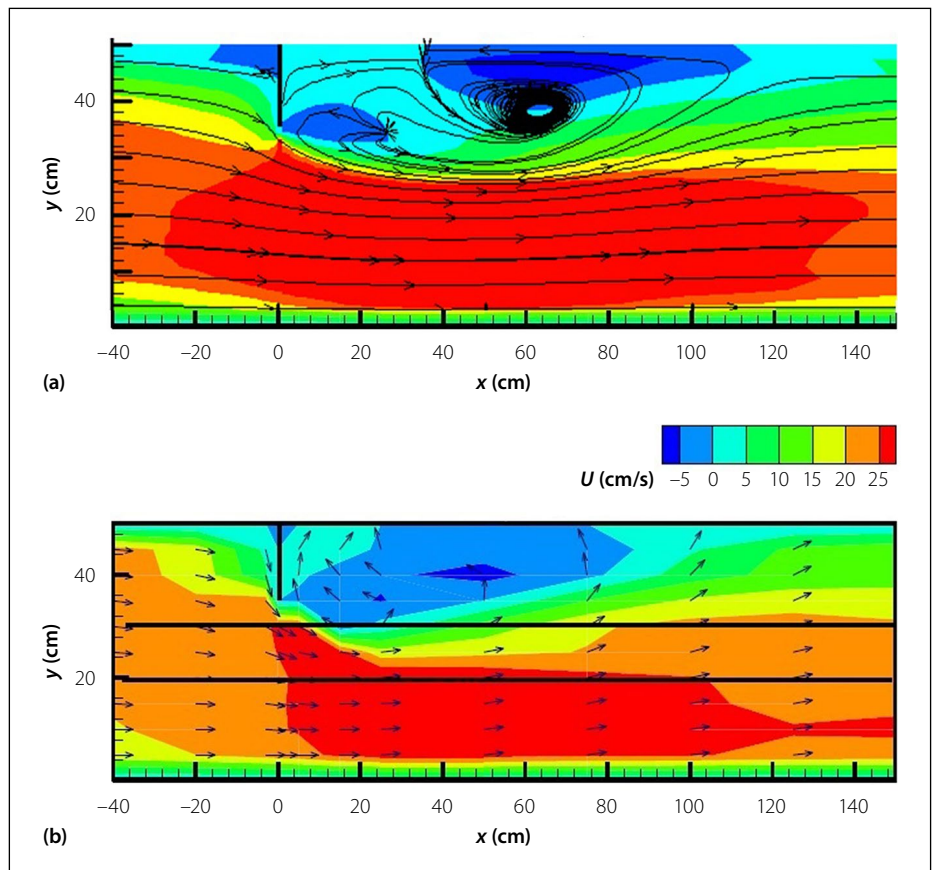


Figure 7 Contour maps of longitudinal velocity (U) for a single groyne on one floodplain: (a) Numerical model (current study), (b) Laboratory model (Ahmed *et al* 2010)

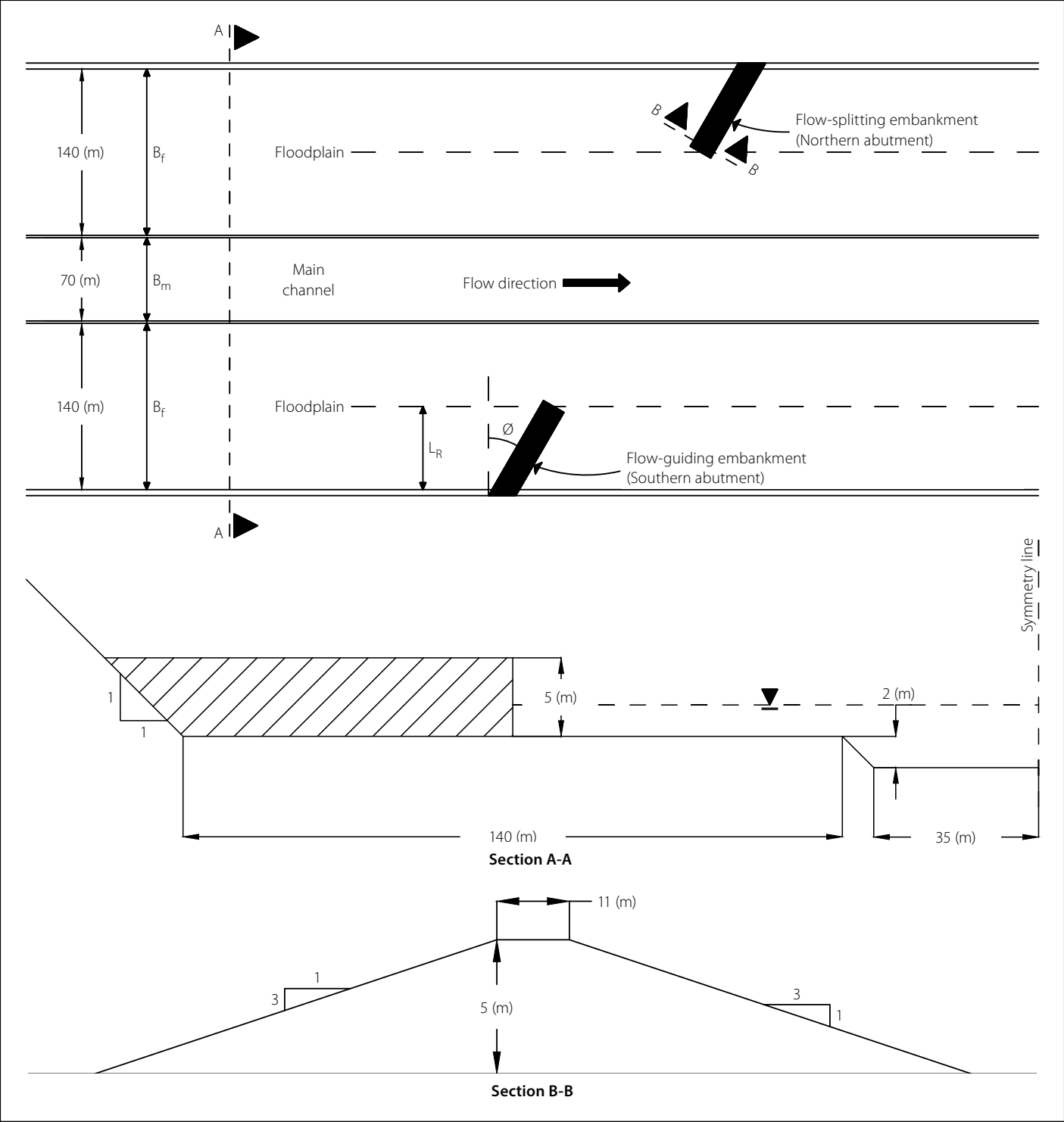


Figure 8 The compound channel dimensions and the skewed bridge crossing

The experiment was conducted using a straight impermeable wooden groyne 0.01 m thick and a relative length (groyne length L_g to the floodplain width) of 0.75 fixed vertically to the main channel centerline and to the flow direction (Ahmed *et al* 2010).

Similar conditions were applied to the developed model for which the contour maps of the longitudinal velocity (U) are shown in Figure 7 where vortices are formed at the groyne downstream. The flow is totally one-way in the laboratory and proposed

numerical models at distances 6.67 and 7.46 times the groyne length, respectively, and the maximum vortex width occurs in them at 1.65 and 1.6 times the groyne length, respectively, concluding that the numerical predictions are accurate and consistent.

Table 2 Bridge skew angles and contraction ratios for different cases

	Case number												
	1	2	3	4	5	6	7	8	9	10	11	12	13
Skew angle (\varnothing)	---	Zero	Zero	Zero	15°	15°	15°	30°	30°	30°	45°	45°	45°
Contraction ratio (L_R/B_f)	---	0.25	0.50	0.75	0.25	0.50	0.75	0.25	0.50	0.75	0.25	0.50	0.75

Results

To perform the current study, some dimensions were assumed for the symmetrical compound channel's cross-section. Floodplain width (B_f) is 140 m, main channel width (B_m) is 70 m, the height difference between the main channel bottom and the

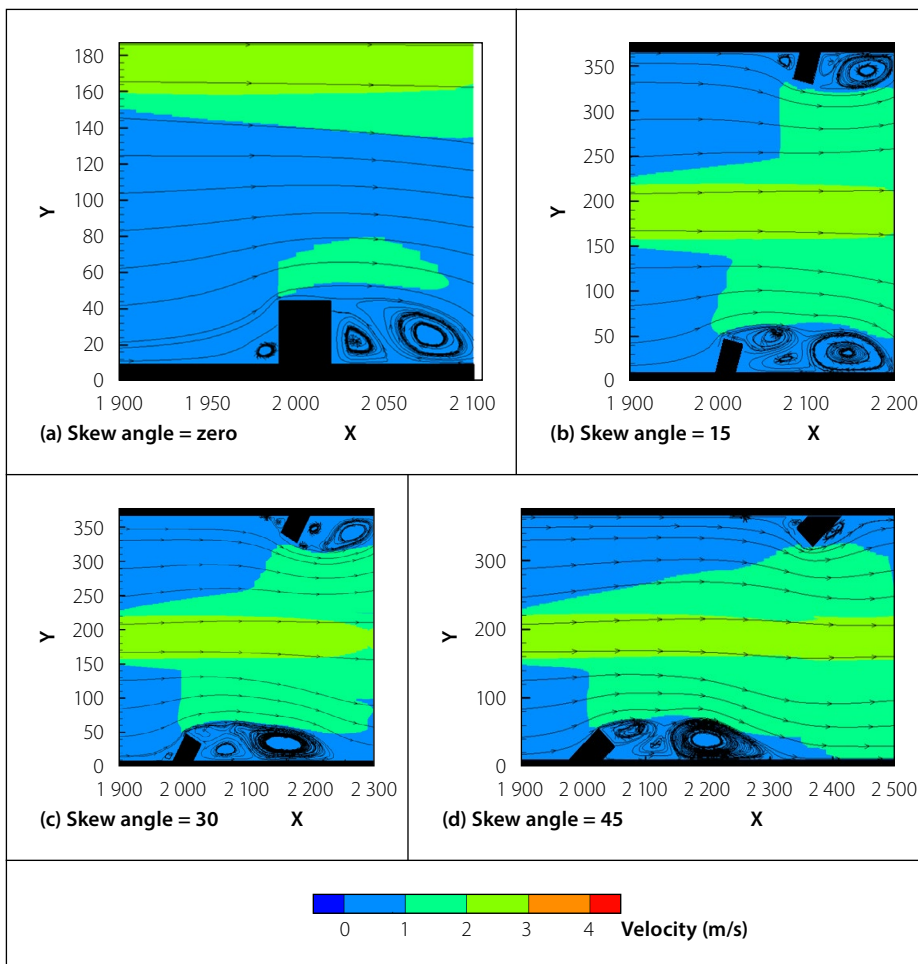


Figure 9 Streamlines and velocities with the contraction ratio of 0.25

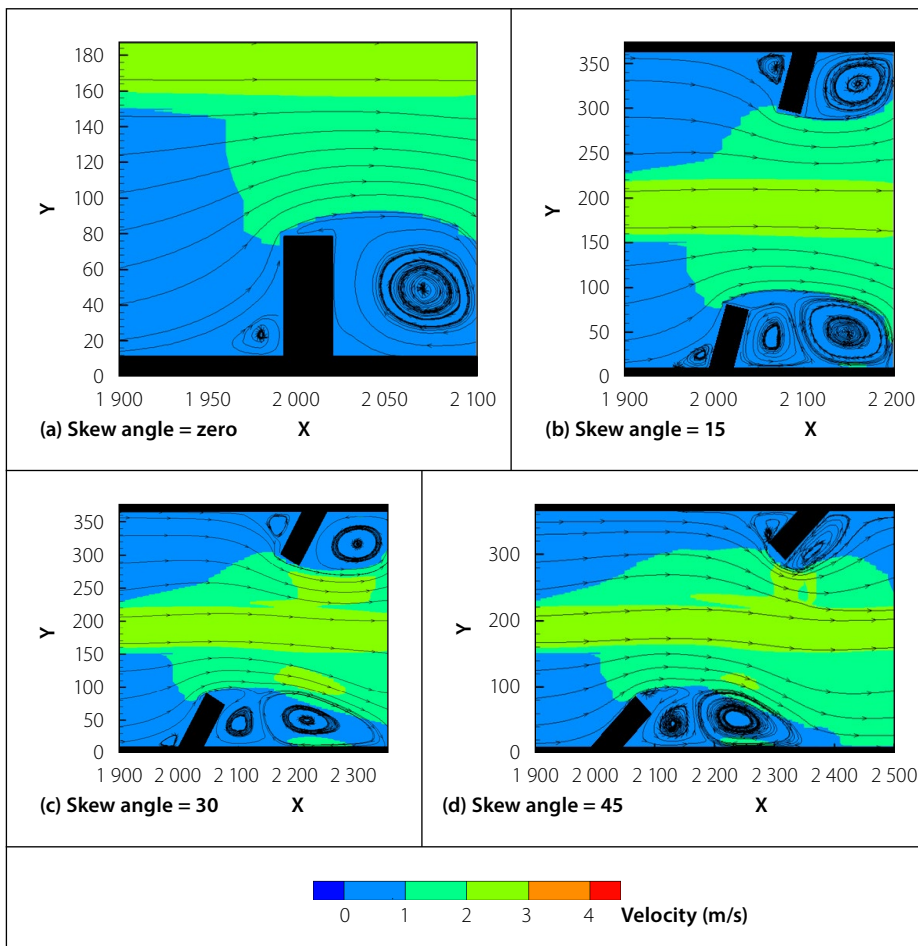


Figure 10 Streamlines and velocities with the contraction ratio of 0.50

floodplain beds is 2 m, the height difference between the floodplain beds and the roadway surface is 5 m, the slope of the sidewalls of the main channel and floodplains is 45 degrees, and the channel's longitudinal slope is 0.0005 (Figure 8 on page 49). The roughness of the main channel bed differs from that of the floodplain; hence, it has been taken as 0.03 for the main channel and 0.05 for the floodplain in all the related simulations. The bridge abutment width in the roadway is 11 m and the slope of the route's side embankments is 3 (horizontal) to 1 (vertical). Also, L_R is the vertical projection of the roadway embankment.

The effects of the roadway contraction ratio (L_R/B_f) and skew angle (θ) on the velocity distribution and bed shear stress have been studied in this research, where the flow-splitting and flow-guiding embankments are expressed, for simplicity, as northern and southern abutments, respectively.

Thirteen cases, with bridge skew angles and contraction ratios as shown in Table 2 (page 49), were considered to study the issue with the numerical model (Case 1 had no bridges) based on the variables shown in Figure 8.

Velocity distribution at the bridge location

Figures 9–11 show the velocity values and streamlines in plan for Cases 2–13 (items for each contraction ratio are shown separately).

These figures show how the flow and velocity are distributed across the bridge. Horizontal vortices, formed mainly behind the abutments, are also visible.

The dimensionless ratio of the maximum velocity in the main channel is shown in Figure 12 where V_{max0} is the maximum velocity in the main channel (Case 1 in Table 2, no-bridge compound channel), V_{max-mc} is the maximum velocity in the main channel area and CR is the contraction ratio.

As shown, when the contraction ratio equals 0.25, an increased skew angle increases the maximum velocity so that at a 45° skew angle (Case 11) the increase is 15% compared to the zero skew angle mode (Case 2). A point worth noting is that, in the main channel, the maximum velocity does not necessarily occur exactly along the bridge axis (parallel to the route); it may increase slightly before or after the bridge location where the southern abutment skew angle seems to direct the flow in the entire opening width along the bridge axis, which reduces the velocity

exactly at the bridge location. A more uniform discharge distribution along the bridge direction reduces the peak velocity in the bridge opening.

When the contraction ratio equals 0.50, an increased skew angle does not change the maximum velocity considerably, because the flows forming behind and in front of the northern and southern abutments do not change significantly, resulting in a more uniform discharge distribution. In this case, the maximum velocity occurs along the bridge axis, except at a skew angle of 45° where it occurs slightly before and after the bridge.

Conditions change somewhat when the contraction ratio equals 0.75. In the main channel with a 15° skew angle, the velocity is 20% more than that with a 0° skew angle, and when the skew angle tends to 45°, the maximum velocity in the main channel nearly equals the no-skew angle mode. In this case, by increasing the skew angle from 15° to 30° and 45°, the length of the road embankment is increased and necessary space is provided to create a recirculation zone upstream of the northern abutment. The formation of this zone guides the flow to the main channel and reduces the peak velocity.

Figure 13 shows the dimensionless ratio of the maximum velocity around the southern abutment (V_{max-sa}) in a bridge-affected area for different contraction ratios. As shown, at all skew angles the maximum velocity around this abutment increases with an increase in the contraction ratio, and the maximum increase (for a fixed skew angle) is related to a 15° skew angle where the difference between contraction ratios of 0.25 and 0.75 is about 164%. The figure also shows that, at all three contraction ratios, an increase in the bridge skew angle reduces the maximum velocity around the southern abutment. The maximum velocity reduction rates at contraction ratios of 0.25, 0.50 and 0.75 are 21%, 46% and 36%, respectively.

According to the flow pattern formed upstream of the southern abutment, it can be concluded that the skew angle causes the southern abutment to act as a guide wall leading the flow uniformly into the main channel at the bridge location which reduces the flow–abutment confrontation.

Around the northern abutment, the situation is more complicated, and Figure 14 shows the dimensionless ratio of the maximum velocity around this abutment (V_{max-na}) for different contraction ratios.

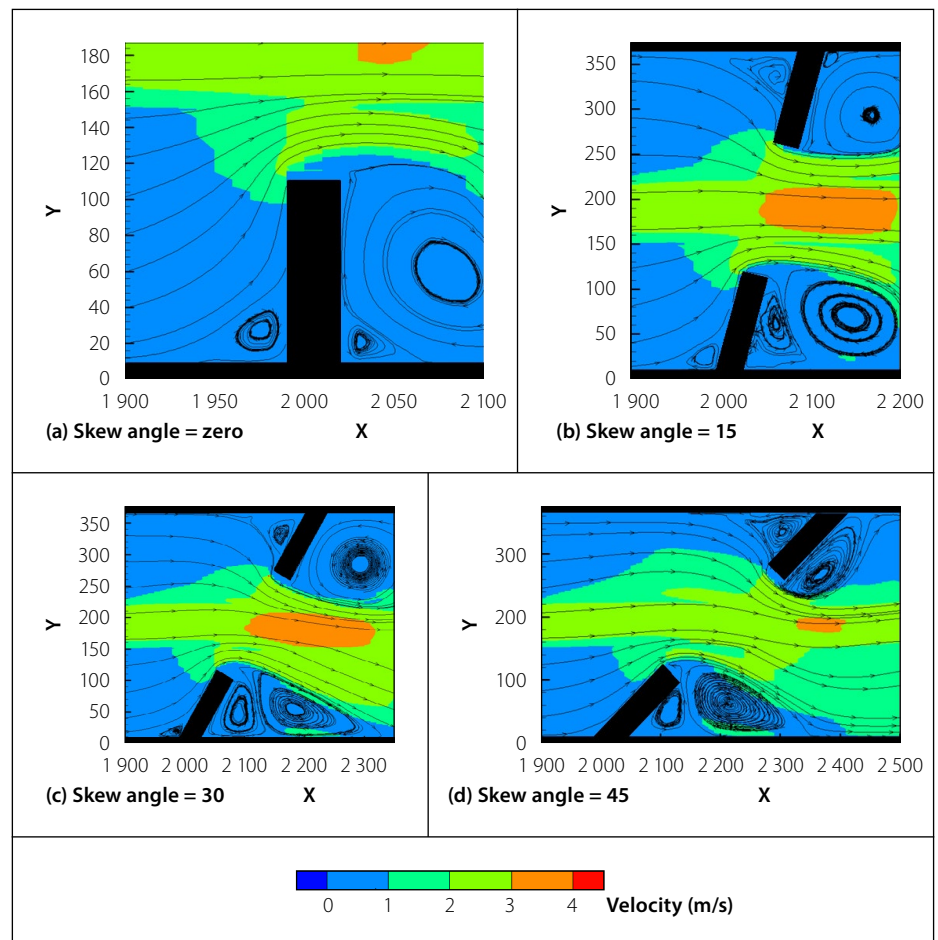


Figure 11 Streamlines and velocities with the contraction ratio of 0.75

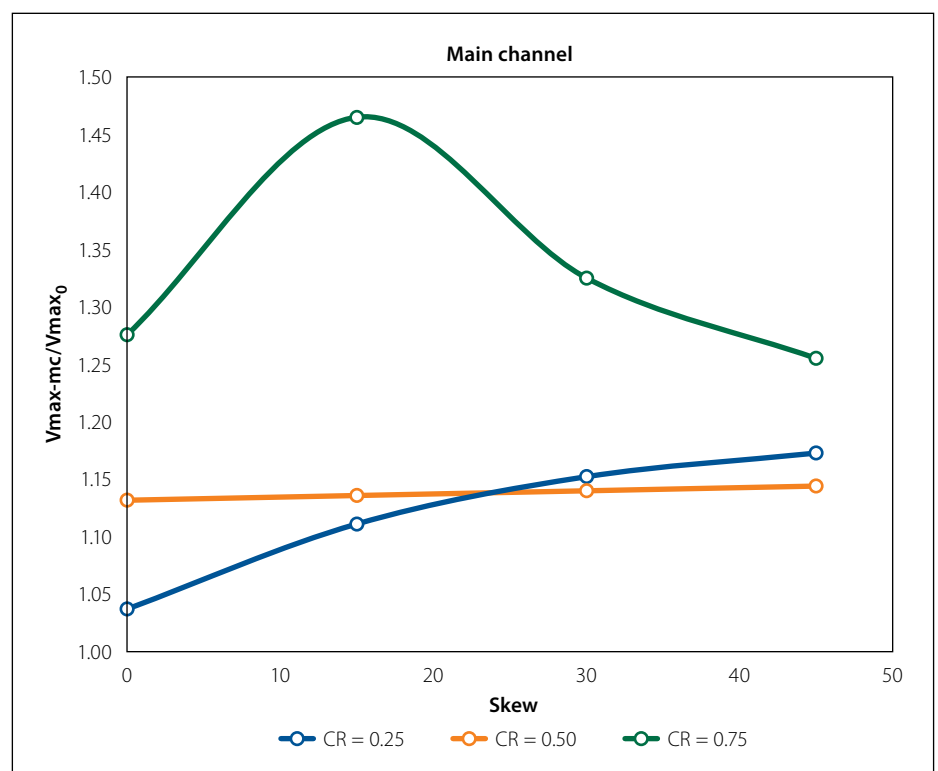


Figure 12 The dimensionless ratio of the maximum velocity in the main channel

As shown, the general trend at contraction ratios of 0.25 and 0.50 is similar for different skew angles; with an increase in the skew angle up to 30°, the maximum velocity around the northern abutment has slightly

reduced, but at a skew angle of 45°, it has increased again (more so for a contraction ratio of 0.25). For skew angles $\leq 30^\circ$, the channel's extended width (along the bridge length) allows the flow to adapt itself to the

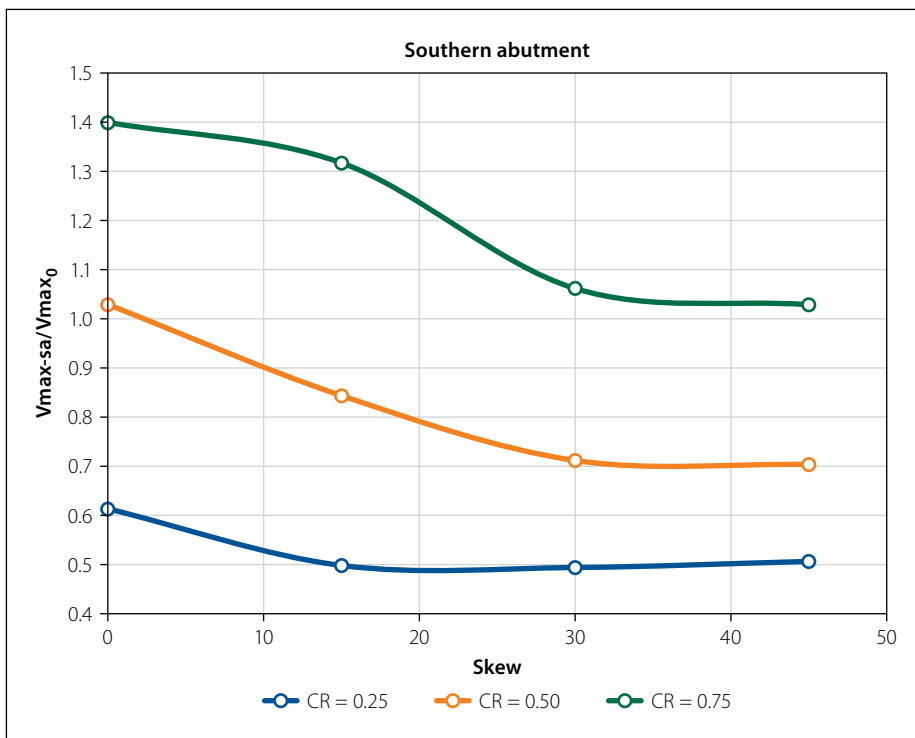


Figure 13 The dimensionless ratio of the maximum velocity around the southern abutment

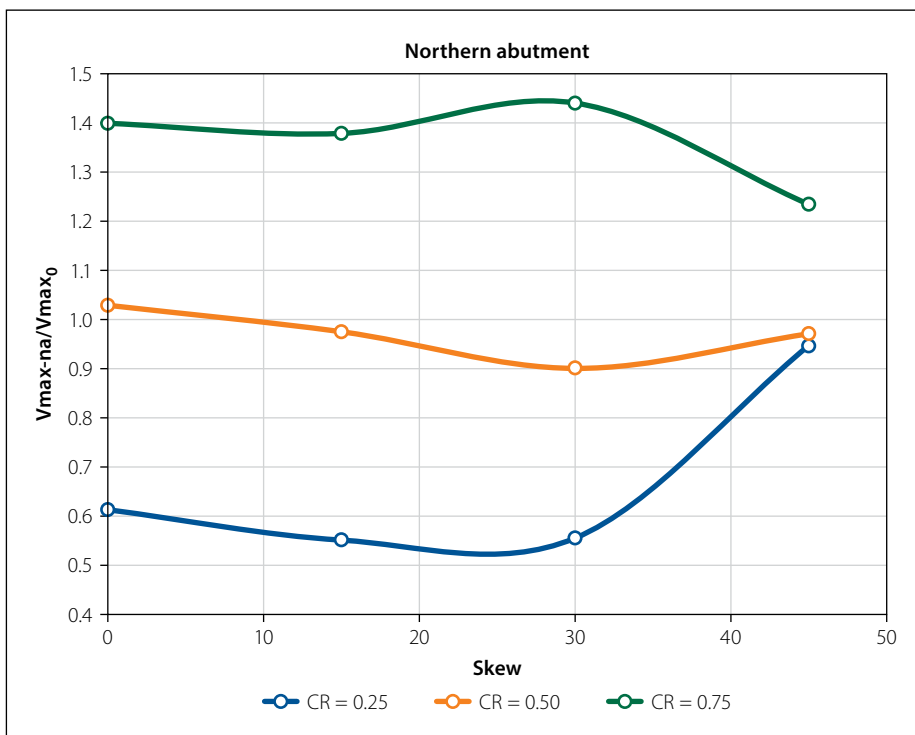


Figure 14 The dimensionless ratio of the maximum velocity around the northern abutment

abutment shape and pass the bridge opening almost evenly, but at a skew angle of 45° the flow–abutment confrontation increases so much that the velocity around the northern abutment is increased. At a contraction ratio of 0.75 where the flood passage gets narrower and the skew angle increases up to 30°, the flow–abutment confrontation causes the maximum velocity to increase slightly around the northern abutment, but at a skew angle of 45° the dead zone formed upstream of this abutment widens and

adjusts the flow circulation so that it passes the bridge opening almost uniformly, causing the maximum velocity to reduce around this abutment.

As mentioned before, the flow around the northern abutment is more complex than that around the southern one because there is a flow–abutment confrontation. Conditions are so complex that defining a specific trend for the flow passing the bridge is very difficult, because the channel and bridge geometry are quite

influential, and each case needs careful study. However, as was predicted, at a skew angle of 0, the increase in the contraction ratio (the decrease in the bridge span) causes the velocity to increase considerably around the northern and southern bridge abutments. The velocity increased by 67% when the contraction ratio increased from 0.25 to 0.50 and by 40% when it increased from 0.50 to 0.75. This velocity increase around the northern abutment can also be observed for the contraction ratio increase for other skew angles (15°, 30° and 45°), with the only difference being at a skew angle of 45° where the percent velocity increase is lower than those of the other cases due to the formation of a dead zone behind the northern abutment.

Another interesting point is that the maximum velocity in the main channel, around the northern abutment and southern abutment, occurs at a contraction ratio of 0.75 and skew angles of 15°, 0° and 30°, respectively. Maximum velocity values are almost equal to each other.

Shear stress at the bridge location

The shear stress parameter is another important flow hydraulic criterion at the bridge abutment. Figure 15 (where S_{max0} is the maximum shear stress in the whole section in the no-roadway embankment mode) shows the nondimensionalised maximum shear stress around the southern abutment (S_{max-sa}) for different contraction ratios.

As shown, for all contraction ratios an increased skew angle reduces the maximum shear stress around the southern abutment. This reduction is moderate at contraction ratios of 0.25 and 0.50, and severe at a contraction ratio of 0.75 (especially at skew angles 0°–30°). For a constant skew angle, as the contraction ratio increases, the maximum shear stress increases, and results confirm this.

As mentioned earlier, conditions were more complex around the northern abutment. Figure 16 shows the dimensionless ratio of the maximum shear stress around the northern abutment (S_{max-na}) for different contraction ratios. When it is 0.25, the maximum shear stress does not change much for skew angles up to 30°, but at a skew angle of 45° the shear stress almost doubles in the northern abutment zone due to the flow–abutment confrontation. At contraction ratios of 0.50 and 0.75, the maximum shear stress is higher than the 0.25 case, but with a different trend. In

these contraction ratios, changes are small at skew angles of zero and 15°, but at 30° the stress is increased considerably due to the flow–abutment confrontation. In both 0.50 and 0.75 contraction ratios, the maximum shear stress decreases sharply due to an increase in the skew angle up to 45° and reaches to about the zero-skew angle case. This is also due to the formation of a dead zone upstream of the abutment that acts like a guide wall causing a uniform flow and a reduction in the maximum shear stress around the northern abutment.

CONCLUSIONS

Determining the flow field and bed shear stress is a necessity for studying the scour of bridges. Since the bridge skew angle and the contraction ratio are the factors that affect the hydraulic performance of river bridges, this paper has studied the flow around bridges in compound channels with different contraction ratios (0.25, 0.50 and 0.75) and various skew angles (0°, 15°, 30° and 45°). Accordingly, a 3D flow model was developed and utilised after its validation by various cases.

When a bridge has a skew angle, one of the abutments directs the flow throughout the opening width, reduces the flow velocity along the bridge axis, and decreases the maximum velocity in the bridge opening due to a more uniform discharge distribution in this direction, especially at small contraction ratios.

When both the bridge skew angle and the contraction ratio are high, the maximum velocity occurs along the bridge axis, instead of in the main channel around the northern abutment where, in general, the flow velocity pattern is complex and determining a similar process for different states is difficult. On the other hand, the maximum velocity around the northern abutment increases for every skew angle, with an increase in the contraction ratio. Again, an increase in the bridge skew angle reduces the maximum velocity around the southern abutment, because the skew angle will cause the roadway embankment to act as a guide wall leading to a uniform flow towards the main channel at the bridge location.

With regard to the bed shear stress, when the contraction ratio does not vary, an increase in the bridge skew angle reduces the maximum shear stress around the southern abutment, and this decrease is more at larger contraction ratios. An increased contraction ratio increases the

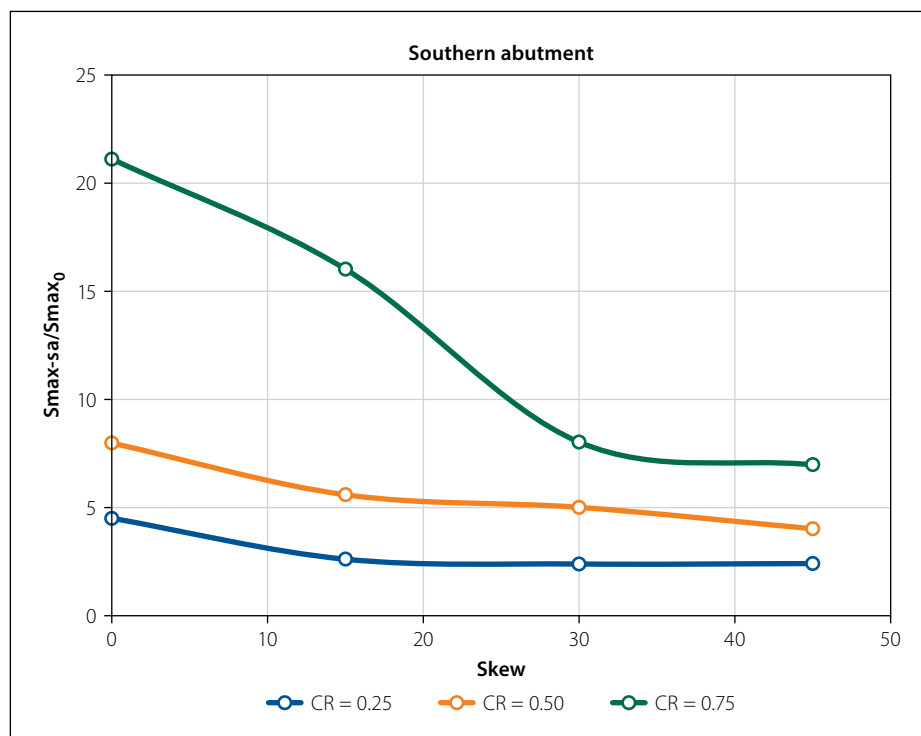


Figure 15 The dimensionless ratio of the maximum shear stress around the southern abutment

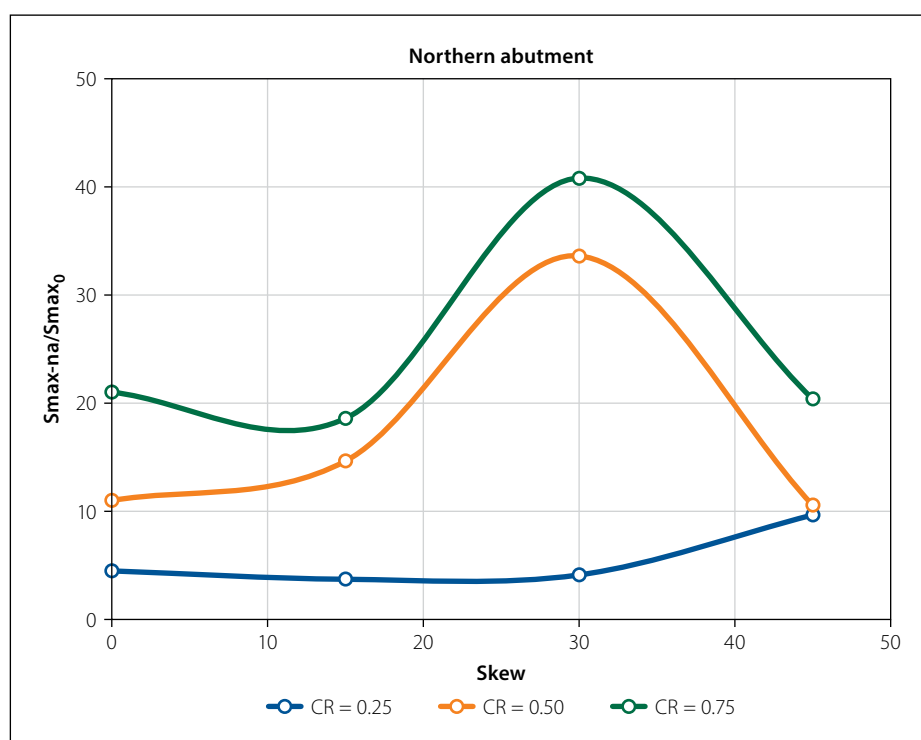


Figure 16 The dimensionless ratio of the maximum shear stress around the northern abutment

maximum shear stress around the mentioned abutments.

Shear stress conditions are also complicated around the northern abutment; at low contraction ratios, the maximum shear stress does not change much at bridge skew angles of up to 30°, but at 45° the shear stress almost doubles in the northern abutment region due to the flow–abutment confrontation. A point worth noting here is the decrease in the maximum shear stress with an increase in the bridge skew

angle when the contraction ratio is high. This, too, can be due to the formation of an upstream flow dead zone that acts like a guide wall causing a uniform flow passage that decreases the maximum shear stress around the northern abutment.

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